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## RESEARCH ARTICLE

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### Key Points:

- Shallow convective cloud ventilation and clear air entrainment collectively shape the clear-to-cloudy boundary layer CO<sub>2</sub>-exchange
- Shallow convective clouds influence the vertical CO<sub>2</sub> distribution until late afternoon, reaching heights twice the boundary layer depth
- Shallow convective clouds shape a strong vertical negative-to-positive CO<sub>2</sub>-H<sub>2</sub>O correlation at shallow cloud-scales

### Supporting Information:

Supporting Information may be found in the online version of this article.

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## Turbulent Exchange of CO<sub>2</sub> in the Lower Tropical Troposphere Across Clear-to-Cloudy Conditions

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**Abstract** This study investigates the roles of clear air entrainment and shallow cloud ventilation, alongside rainforest CO<sub>2</sub>-assimilation, in the turbulent exchange of CO<sub>2</sub> within the lower tropical troposphere under clear-to-cloudy conditions. Constrained by comprehensive observations from the CloudRoots-Amazon22 campaign, spanning leaf stomatal to upper atmosphere, we design and evaluate a representative shallow convective numerical experiment with the turbulence-resolving Dutch Atmospheric Large Eddy Simulation model, incorporating a bulk rainforest representation. We assess contributions from the rainforest, clouds, and environment through the vertically integrated, domain-averaged CO<sub>2</sub> budget by comparing simulations with and without the dynamic effects of clouds. Our findings reveal three distinct diurnal regimes named: entrainment-diluting, cloud-ventilation-and-entrainment, and CO<sub>2</sub>-assimilation. Shallow convective clouds (~23%), clear air entrainment (~21%), and rainforest CO<sub>2</sub>-assimilation (~56%) collectively influence the diurnal evolution and vertical exchange of CO<sub>2</sub> within the clear-to-cloudy boundary layer, with their relative importance varying per diurnal regime. In the absence of clouds, ventilation ceases and CO<sub>2</sub> exchange is driven solely by entrainment and CO<sub>2</sub>-assimilation, resulting in a 20%–25% reduction in mixing effectiveness. In the vertical, shallow clouds ventilate CO<sub>2</sub> to heights reaching twice the boundary layer depth, significantly affecting the vertical distribution until late afternoon. Analysis of the correlation between CO<sub>2</sub> and H<sub>2</sub>O shows that shallow convective clouds organize the turbulent exchange at shallow cloud-scales, shaping a vertical pattern of negative to positive CO<sub>2</sub>-H<sub>2</sub>O correlation from the roughness sublayer into the cloud layer. These findings highlight key processes crucial for accurately representing the lower tropical tropospheric CO<sub>2</sub> budget across clear-to-cloudy conditions.

**Plain Language Summary** Understanding climate change requires a full picture of how carbon dioxide (CO<sub>2</sub>) behaves in the atmosphere, particularly over tropical rainforests such as the Amazon, where interactions between the atmosphere and rainforest are strongest. To capture this behavior, we must study the system from the rainforest to the upper atmosphere. Although progress has been made in understanding how rainforests exchange CO<sub>2</sub>, major uncertainties remain; especially when clouds are present. Clouds significantly affect how CO<sub>2</sub> moves in the lower atmosphere compared to clear conditions. To investigate this further, we combined detailed observations from the 2022 CloudRoots-Amazon22 campaign with a high-resolution model that simulates turbulence and clouds. After a thorough validation of the simulation, our results show that clouds and the clear air environment actively shape the behavior of CO<sub>2</sub>. During daytime, we identified three distinct periods in the daily cycle that describe CO<sub>2</sub> movement when clouds start to form in an initially clear sky. Vertically, clouds transport CO<sub>2</sub> from near the surface up to heights of 3.5 km. Moreover, clouds not only lift CO<sub>2</sub> but also organize how it moves upward. This research highlights key processes that are crucial for accurately understanding and representing CO<sub>2</sub> dynamics in the atmosphere above tropical rainforests.

## 1. Introduction

Obtaining a full understanding of anthropogenic climate change necessitates a thorough understanding of the diurnal (daytime, between sunrise and sunset) and vertical variability of atmospheric carbon dioxide (CO<sub>2</sub>) above tropical rainforests such as the Amazon, the largest terrestrial carbon sinks on our planet (Anav et al., 2015; Friedlingstein et al., 2023). Although significant progress has been made in quantifying rainforest carbon exchange and its response to extreme weather events such as droughts and heatwaves (e.g., Alden et al., 2016; Gatti

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et al., 2014; van der Laan-Luijkx et al., 2015), uncertainties remain in estimates of these surface exchange fluxes, largely due to limited observations and limitations in accurately representing the atmospheric processes involved in the CO<sub>2</sub> exchange (Bastos et al., 2020; Botía et al., 2022).

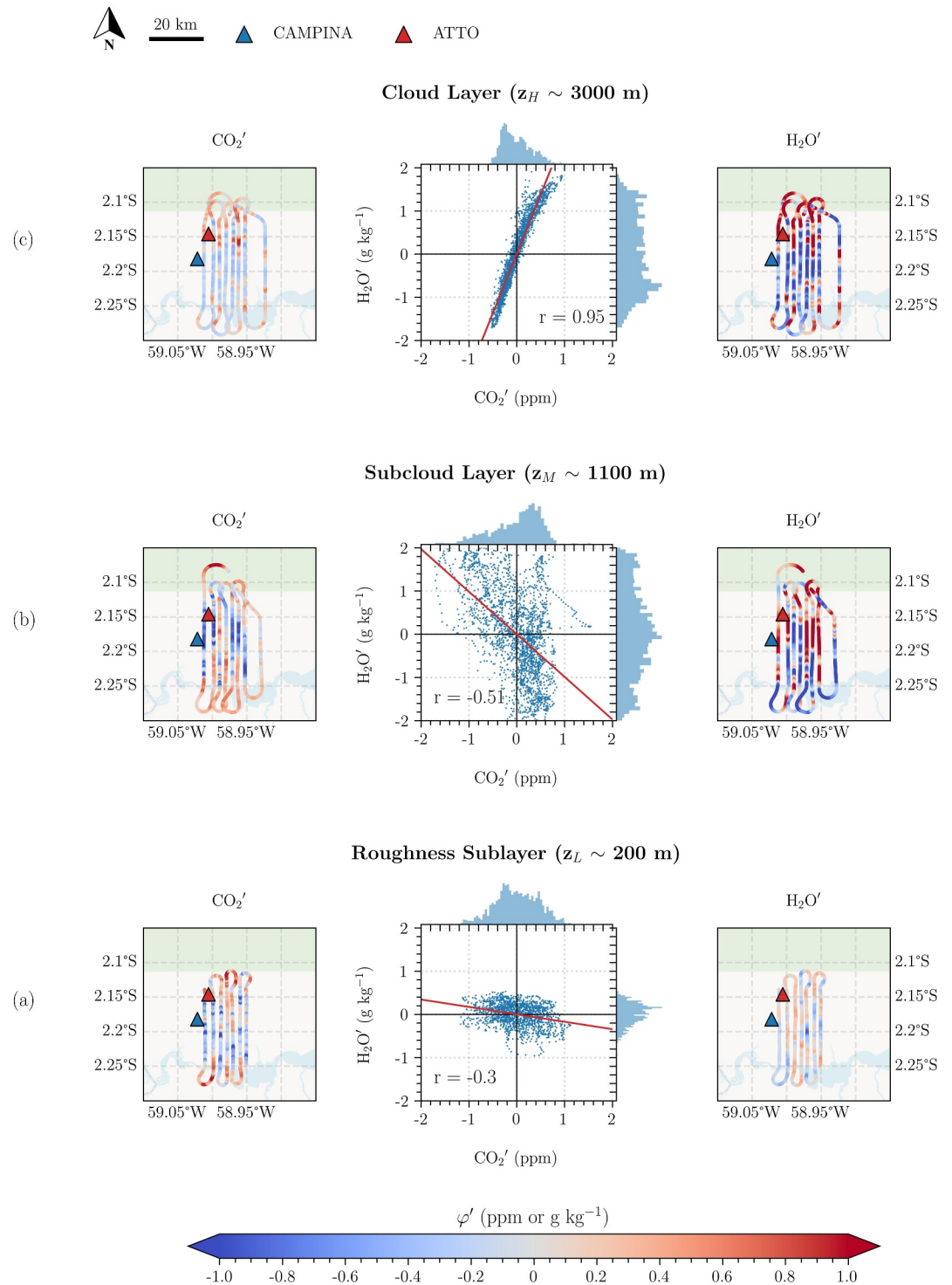
Among the atmospheric processes governing CO<sub>2</sub> exchange, moist convection remains particularly unexplored. Previous work has emphasized the importance of biospheric to atmospheric fluxes of photosynthesis and respiration (e.g., Fitzjarrald et al., 1990; Garstang & Fitzjarrald, 1999), dry convection (e.g., Hooft et al., 2019; Pino et al., 2012), and horizontal advection (e.g., Botía et al., 2022; Pino et al., 2013). However, the contribution of moist convective processes to the CO<sub>2</sub> budget and subsequently the errors arising from misrepresenting these processes, remain elusive. Particular attention should be given to the role of shallow moist convection, which is a characteristic feature above tropical ecosystems such as the Amazon during the dry season (Giangrande et al., 2020).

During the morning, the clear boundary layer above the dry season rainforest surface typically transitions to a layer topped by small, active, and non-precipitating shallow convective clouds (<2 km in depth) (Giangrande et al., 2020; Henkes et al., 2021). These active clouds are closely coupled to the forest (Gentine et al., 2019). Hence, shallow moist convection is thought to significantly influence and connect the diurnal cycles of energy, water, and atmospheric CO<sub>2</sub> in the lower troposphere (Gentine et al., 2019; Sikma & Vilà-Guerau de Arellano, 2019; Vilà-Guerau de Arellano et al., 2020). To place such ideas on more quantitative grounds, we will study the contributions of dry and moist convection, alongside CO<sub>2</sub>-assimilation by the Amazon rainforest, in shaping the diurnal and vertical variability of atmospheric CO<sub>2</sub> in the lower tropical troposphere.

To more specifically motivate how shallow moist convection might influence the vertical variability of CO<sub>2</sub> over the Amazon, we present aircraft observations from the recent CloudRoots-Amazon22 measurement campaign (Vilà-Guerau de Arellano et al., 2024) (Figure 1). During this campaign, held at the Amazon Tall Tower Observatory (ATTO) and Campina measurement site (Andreae et al., 2015), an aircraft equipped with Picarro gas analyzers, flew a vertical profile combined with three horizontal rasters with a distance of 20 km in the early afternoon (13 hr Local Time (LT) = UTC−4 hr) on the 18<sup>th</sup> of August 2022, where non-precipitating shallow convective clouds formed from the late morning onward. The horizontal levels were distributed throughout the lower troposphere: in the roughness sublayer (low (L),  $z_L$  at ~200 m), subcloud layer (medium (M),  $z_M$  at ~1,100 m), and cloud layer (high (H),  $z_H$  ~ 3,000 m). For simplicity the height of the subcloud layer is taken synonymous to the well-mixed boundary layer height when shallow convective clouds are present (well-mixed boundary layer height ( $z_i$ ) at ~1,500 m as measured by a ceilometer). During the flight, the aircraft measured mole fractions of  $\varphi \in [\text{CO}_2, \text{H}_2\text{O}]$  in the environment; the aircraft did not fly through the clouds themselves.

In Figure 1 we visualize the linearly detrended correlation between fluctuations ( $\varphi'$ ) in H<sub>2</sub>O and CO<sub>2</sub> around spatial raster-averages ( $\overline{\varphi}$ ). The figure shows that CO<sub>2</sub>' and H<sub>2</sub>O' are correlated at each height and that this correlation changes with height. In the roughness sublayer, a region that is still influenced by the rainforest, CO<sub>2</sub> and H<sub>2</sub>O are negatively correlated (Pearson correlation coefficient ( $r$ ) = −0.3), with larger CO<sub>2</sub>' compared to H<sub>2</sub>O'. Higher up in the cloud layer (3,000 m), CO<sub>2</sub> and H<sub>2</sub>O are strongly positively correlated ( $r$  = 0.95), as variability in CO<sub>2</sub> decreases, whereas variability in H<sub>2</sub>O increases. In between, in the upper part of the subcloud layer ( $z_M$ ), the correlation is also negative ( $r$  = −0.51), although there is a considerable variability in both H<sub>2</sub>O and CO<sub>2</sub>.

The vertical transition from negative to positive CO<sub>2</sub>-H<sub>2</sub>O correlation up to  $z_i$  was previously noted by Vilà-Guerau de Arellano et al. (2004) above grassland at Cabauw (The Netherlands), under non-cloudy conditions. They concluded that near the surface, CO<sub>2</sub> mole fractions decrease through descending motions induced by photosynthesis (CO<sub>2</sub>-assimilation), whereas ascending convective thermals originating from the surface are moistened from below through evapotranspiration, creating a strong negative correlation in the roughness sublayer. Near  $z_i$ , the CO<sub>2</sub> and moisture in an ascending thermal originating from the surface will have relaxed toward typical boundary layer CO<sub>2</sub> mole fractions through lateral mixing. Since the ambient CO<sub>2</sub> and moisture generally reduce with height around  $z_i$ , a thermal with bulk-boundary layer moisture and CO<sub>2</sub> will possess both a CO<sub>2</sub> and moisture excess. Conversely, compensating descending motions from the free troposphere will contain both less CO<sub>2</sub> and moisture. Hence, when both the CO<sub>2</sub> and moisture gradients with height are negative, as is often the case at  $z_i$  and in the lower free troposphere, we expect both clear and cloudy circulations to generate positive CO<sub>2</sub>-moisture correlations around  $z_i$ . However, in the observations displayed in Figure 1, a positive correlation is found



**Figure 1.** Observed linearly detrended correlation between fluctuations ( $\phi'$ ) in  $\text{CO}_2$  and  $\text{H}_2\text{O}$  around spatial raster-averages ( $\bar{\phi}$ ) and the corresponding Pearson correlation coefficient ( $r$ ) at 13 LT across three heights: (a)  $z_L$  (roughness sublayer at  $\sim 200$  m), (b)  $z_M$  (subcloud layer at  $\sim 1,100$  m) and (c)  $z_H$  (cloud layer at  $\sim 3,000$  m). Note that for simplicity the height of the subcloud layer is taken synonymous to the well-mixed boundary layer height when shallow convective clouds are present (well-mixed boundary layer height ( $z_i$ ) at  $\sim 1,500$  m as measured by a ceilometer). Basemap: OpenStreetMap, with water (blue), terra firm (brown), and rainforest (green). Triangular markers indicate the location of the Amazon Tall Tower Observatory (ATTO) and Campina measurement site.

far above  $z_i$ , in the shallow cloud layer, suggesting that shallow convective clouds may play a role in extending this vertical positive correlation.

Shallow convective clouds can transport (“vent”) volumes of air from the subcloud layer into the cloud layer and across the lower troposphere via upward mass fluxes ( $M$ ) (Nicholls & Lemone, 1980; Ouwersloot et al., 2013; van Stratum et al., 2014). This venting typically dries out and limits the growth of  $z_i$ , whereas also influencing the atmospheric residence time of transported gases (Bardakov et al., 2022; Lu et al., 2000). Simultaneously, shallow convective clouds facilitate downward transport to the surface (“pump”), both locally, through subsiding shells around the clouds and across broader regions via compensating subsidence (Jonker et al., 2008; McMichael et al., 2020; Rodts et al., 2003). Consequently, these venting and pumping characteristics raise the question whether the strong positive  $\text{CO}_2$ - $\text{H}_2\text{O}$  correlation within the cloud layer reflects upward  $\text{CO}_2$  transport via moist cloudy mass fluxes and downward transport through clear skies, and whether this correlated signal is present and remains throughout the day when we move from clear (dry convective) to cloudy (moist convective) conditions. More concretely, Figure 1 motivates us to formulate three research questions:

1. How do cloud ventilation by shallow convective clouds, clear air entrainment, and rainforest  $\text{CO}_2$ -assimilation collectively regulate the turbulent exchange of atmospheric boundary layer  $\text{CO}_2$  across clear-to-cloudy conditions?
2. How does cloud ventilation by shallow convective clouds influence the vertical distribution of  $\text{CO}_2$  in the lower tropical troposphere?
3. What role does shallow moist convection play in shaping the vertical correlation pattern between  $\text{CO}_2$  and water vapor, as illustrated in Figure 1?

To address these questions, we combine observations from the CloudRoots-Amazon22 campaign to design and analyze a Large Eddy Simulation (LES) experiment for a typical non-precipitating shallow convective day in the Amazonian dry season. We run two simulations: one standard with clouds and one mechanism-denial experiment without clouds, where condensation is prevented from affecting model dynamics. The observations and details of the numerical experiment are introduced in Section 2. We first validate the simulation in Section 3.1. Next, we examine and discuss the role entrainment (dry convection) and mass fluxes (moist convection), alongside rainforest  $\text{CO}_2$ -assimilation, in the  $\text{CO}_2$  budgets of the clear-to-cloudy boundary layer (Section 3.2) and lower troposphere (Section 3.3). Furthermore, we examine whether turbulence and shallow convective clouds organize vertical exchange of  $\text{CO}_2$  as represented in Figure 1 (Section 3.4). Lastly, we summarize the key findings in Section 4.

## 2. Material and Methods

### 2.1. CloudRoots-Amazon22 Campaign

The CloudRoots-Amazon22 campaign (Vilà-Guerau de Arellano et al., 2024) was conducted between the 8<sup>th</sup> to 21<sup>st</sup> August 2022 (dry season) at the Amazon Tall Tower Observatory (ATTO) (lat:  $-2.1458^\circ$ , lon:  $-59.0055^\circ$ ) and Campina (lat:  $-2.1819^\circ$ , lon:  $-59.0217^\circ$ ) measurement sites located in central Amazon (Andreae et al., 2015). During the campaign, aimed at studying rainforest-atmosphere exchange across the ~40 m tall canopy and lower troposphere, continuous measurements of soil properties and atmospheric state variables (including  $\text{CO}_2$  mole fractions), along with their turbulent fluxes, were conducted on and around two towers of 323 and 81 m height (Vilà-Guerau de Arellano et al., 2024). The vertical domain was expanded by releasing daily radiosondes at a three-hour interval (6, 9, 12, 15, and 18 LT (LT = UTC−4 hr)). Additionally, on the 18<sup>th</sup> of August 2022, an aircraft measured vertical profiles (up to 5 km) of both  $\text{CO}_2$  and humidity, and flew the 20 km rasters presented in Figure 1. In addition to the afternoon (13 LT) flight presented in Figure 1, aircraft observations were also collected during the morning (9 LT).

To initialize and evaluate our LES experiment, we use the following observations. First, we combine the meteorological state variables and  $\text{CO}_2$  mole fractions from the towers, radiosondes, and aircraft into connected vertical profiles from within the canopy to the mid-troposphere according to their availability. In addition, we infer the boundary layer height ( $z_i$ ) from the observed backscattering of a permanently installed ceilometer at Campina and from the available radiosondes. Regarding the radiosondes, we follow well-mixed theory and  $z_i$  is taken as the height where fits of a constant well-mixed potential temperature value, derived from the lowest levels of the connected vertical profiles, intersects a linear fit of the free-tropospheric lapse rate of potential temperature

(e.g., Lilly, 1968; Pino et al., 2006). Additionally, we evaluate simulated clouds against reflectivity, type, base, thickness and top retrieved from a MIRA Doppler cloud radar. The ceilometer and MIRA radar did not operate on the 17<sup>th</sup> of August 2022, and thus data for this day is missing.

Based on the observations, visual inspection of available time-lapse imagery from an upward looking camera at ATTO, and previous characterization studies (e.g., Henkes et al., 2021), data from the CloudRoots-Amazon22 campaign were categorized into two convective regimes based on the type of clouds that developed throughout the day: shallow convective (6 days, 9, 10, 11, 15, 17, and 18 of August 2022) and shallow-to-deep convective (4 days, 12, 13, 14, and 16 of August 2022) (Vilà-Guerau de Arellano et al., 2024). This study focuses solely on observations from the shallow convective regime.

## 2.2. Numerical Simulation

Simulations of meteorological state variables and CO<sub>2</sub> are performed with the Dutch Atmospheric Large Eddy Simulation (DALES) (Heus et al., 2010; Ouwersloot et al., 2017), version 4.4. The simulation design builds on the Green Ocean Amazon 2014/15 case studied by Vilà-Guerau de Arellano et al. (2020). Its domain spans 19 × 19 km horizontally ( $\Delta x = \Delta y = 53$  m) to match the aircraft raster size, and 5 km vertically ( $\Delta z = 20$  m). The model is run for 12 hr ( $\Delta t = 1$  s) between 6 and 18 LT and is horizontally constrained by doubly periodic boundaries, where the first hour is considered spin-up. In the vertical, the bottom surface is forced by a coupled land surface model, whereas the top of the domain contains a sponge layer (Heus et al., 2010). The sponge layer is a relaxation toward the top of the domain in which we dampen the fluctuations of state variables and CO<sub>2</sub> generated by clouds and turbulence in the upper part of the domain around their mean state to prevent numerical instability. For our simulation, we consider the upper 25% (between 3.75 and 5 km) of the domain for the sponge layer, which is at a sufficient height from the shallow convective cloud tops (at a height of ~3 km, Figure 2c) in order to not intervene with their development.

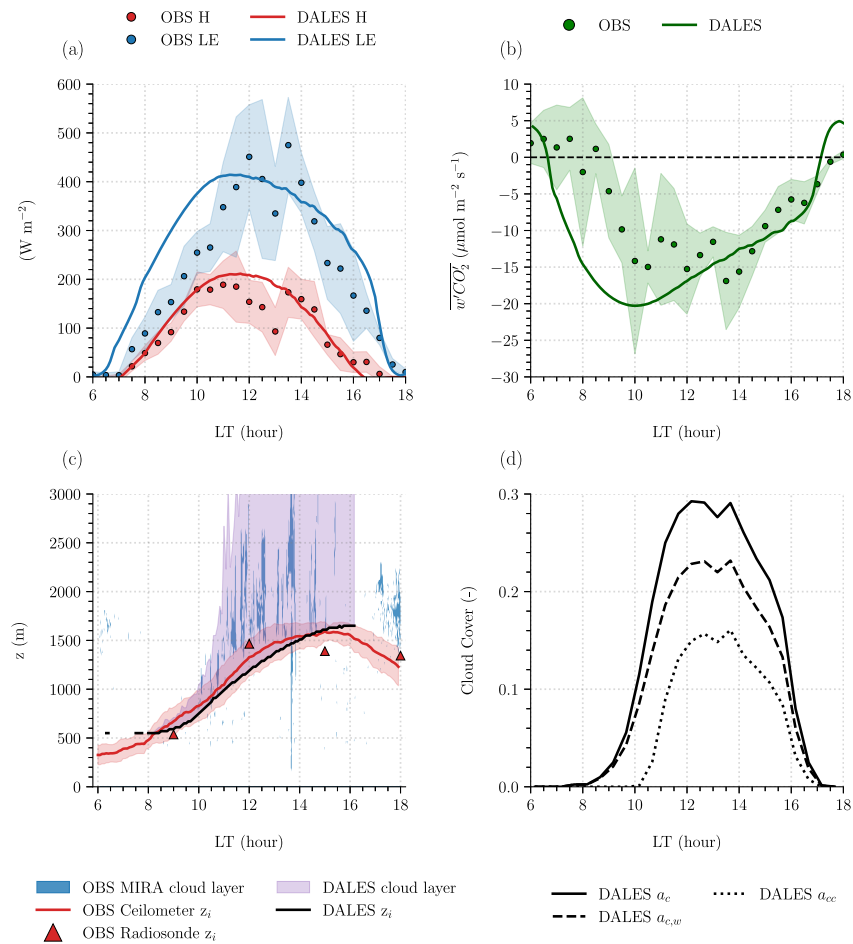
The land surface model is a CO<sub>2</sub>-assimilation (A) represented by stomatal conductance ( $g_s$ ) A- $g_s$  model (Pedruzzi-Bagazgoitia et al., 2017; Ronda et al., 2001), constrained to the phenological observations taken during CloudRoots-Amazon22 (González-Armas et al., 2025; Vilà-Guerau de Arellano et al., 2024). The  $g_s$  represents the stomatal aperture or opening of the stomata, controlling the in- or outflow of CO<sub>2</sub> and H<sub>2</sub>O in the A- $g_s$  scheme. The stomatal aperture depends on environmental conditions such as photosynthetically active radiation, temperature and vapor pressure deficit, which feed back into the surface energy and CO<sub>2</sub> fluxes (Ronda et al., 2001). For the A- $g_s$  scheme a single sunlit big-leaf representation is used for each grid cell as a simplified representation of all individual leaves of the rainforest in the horizontally homogeneous domain. The A- $g_s$  scheme is upscaled and incorporated into a bulk (vertically integrated) rainforest canopy representation. Additionally, the daytime respiration is fixed at 5.18  $\mu\text{mol m}^{-2} \text{s}^{-1}$  as observations displayed near-constant respiration at this rate (Figure S1 in Supporting Information S1). For radiation, we use a one-dimensional Rapid Radiative Transfer Model for Global circulation (RRTMG) (Iacono et al., 2008), calculating partitioning of radiation into direct and diffuse radiation (including photosynthetically active radiation) at each column of the model. Consequently, cloud shadows are projected directly below the cloud.

Although using a sunlit big-leaf approach and the one-dimensional RRTMG scheme may lead to misrepresentations of the meteorology and subsequent surface fluxes such as CO<sub>2</sub>-assimilation, our DALES simulation shows to reproduce the radiation, boundary layer development, meteorological state variables, and surface fluxes satisfactorily compared to the CloudRoots-Amazon22 observations (Figures 2 and 3, Section 3.1). Therefore, we expect any discrepancies arising from these model settings to be negligible for assessing the diurnal and vertical variability of CO<sub>2</sub> in the clear-to-cloudy lower troposphere.

Within the LES domain, CO<sub>2</sub> is advected with the resolved and unresolved flow as an inert tracer. Model grid boxes are either entirely cloudy or cloud-free, and condensation affects the local thermodynamics by instantaneous saturation adjustment (“all or nothing” cloud adjustment scheme (Heus et al., 2010)). All rain microphysics are disabled, as we investigate non-precipitating shallow convective clouds.

We initialize the simulation with vertical profiles of potential temperature ( $\theta$ ), specific humidity ( $q$ ), and horizontal winds ( $u$ ,  $v$ ) averaged (aggregated) over the six observed shallow convective days of CloudRoots-Amazon22 (Vilà-Guerau de Arellano et al., 2024), at 6 LT. For CO<sub>2</sub>, an initial vertical profile is constructed from the morning vertical flight profile (9 LT) above  $z_i$  (vertical (initial) profiles of  $\theta$ ,  $q$ , CO<sub>2</sub>,  $u$ , and  $v$  are

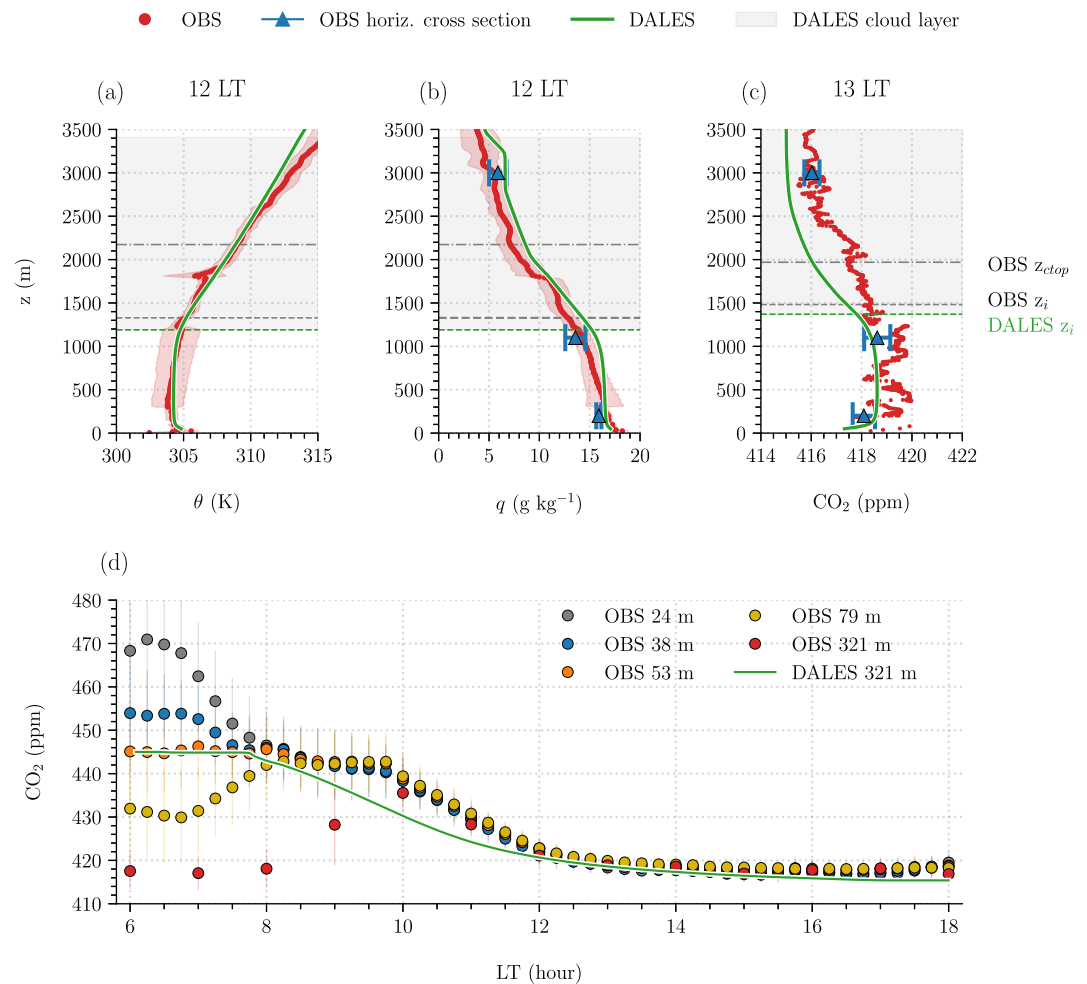




**Figure 2.** Observed (markers) and simulated (solid line) 6-day aggregated diurnal evolution of (a) rainforest surface fluxes—sensible (H) and latent heat flux (LE), (b)  $\text{CO}_2$  surface flux ( $\overline{w'\text{CO}_2}$ )—at 10 m above canopy top ( $z = 50$  m), along with (c) 6-day aggregated diurnal development of the clear-to-cloudy boundary layer ( $z_i$ ), cloud layer and cloud onset, and (d) the diurnal evolution of the simulated mean cloud fraction, split into the three criteria: cloud-average cloud fraction ( $a_c$ , where  $q_l > 0 \text{ g kg}^{-1}$ ), upward-moving cloud-average cloud fraction ( $a_{c,w}$ , where  $q_l > 0 \text{ g kg}^{-1}$ ,  $w' > 0 \text{ m s}^{-1}$ ) and core-average cloud fraction ( $a_{cc}$ , where  $q_l > 0 \text{ g kg}^{-1}$ ,  $w' > 0 \text{ m s}^{-1}$  and  $\theta'_e > 0 \text{ K}$ ). Markers indicate the observed 6-day aggregated diurnal behavior for days on which shallow convective clouds were observed, whereas the shaded areas represent the spread of one standard deviation during the six observed shallow convective days. For DALES,  $z_i$  is set equal to the height of the minimum buoyancy flux.

presented in Supplementary Figures S2 and S3 in Supporting Information S1). Below  $z_i$ , tower  $\text{CO}_2$  data suggest that increased  $\text{CO}_2$  mole fractions ( $>450$  ppm) from respiration in the early morning (6–10 LT), which are transported upward (flushed-up) from near the surface as turbulence initiates during the morning transition (Culf et al., 1997; Dupont et al., 2024; González-Armas et al., 2025) (Figure 3d). Our bulk canopy approach cannot directly represent these in-canopy motions which create this flush-up. The A- $g_s$  scheme misses a method for representing the morning transition in which we move from nighttime respiration to the daytime assimilation. Consequently, assimilation activates immediately after the spin up of the simulation. To ensure that we do simulate realistic conditions by the time the flush-up is completed, we manually enhance the inversion strength from the 9 LT profile in our initial profile based on the tower (24, 38, 53, 79, and 321 m) observations, matching near-canopy (53 m) observations and attaining a realistic well-mixed  $\text{CO}_2$  mole fraction in the course of the day. The selected values of the phenological and model parameters, model settings, and initial profiles used for the simulations are accessible via the supporting code and data repository (de Feiter et al., 2025).

To isolate the role of shallow convective clouds, we impose no domain-averaged horizontal advection and vertical advection, assuming a zero large-scale vertical velocity. We also conduct two simulations; one with and one



**Figure 3.** Observed (markers) and simulated (solid line) vertical profiles at 12 and 13 LT: (a) potential temperature ( $\theta$ ) and (b) specific humidity ( $q$ ) from daily radiosonde observations and (c) CO<sub>2</sub> from aircraft observations on the 18<sup>th</sup> of August 2022. Markers indicate the mean profile, whereas the shaded areas represent the spread of one standard deviation during the six observed shallow convective days. The observed vertical profiles combine tower and radiosonde data, explaining the lower standard deviation in the lower 300 m. Blue horizontal markers represent the horizontal cross sections as measured by the aircraft at 13 LT on the 18<sup>th</sup> of August 2022, horizontal dashed or dashed-dotted lines represent the observed (gray lines) and simulated (green line)  $z_i$  and average cloud top at ATTO ( $z_{ctop}$ ), respectively. Additionally, (d) shows the diurnal evolution of CO<sub>2</sub>, observed (marker) and simulated (green line) mole fraction at 24, 38, 53, 79, and 321 m between 6 and 18 LT. Markers indicate the mean observations, whereas the error bars denote the spread of one standard deviation during the six observed shallow convective days.

without dynamically active shallow convective clouds as a mechanism-denial experiment, where condensation of water vapor in clouds does not influence the radiation and thermodynamic calculations; the latent heat of vaporization is set to zero. These simulations are referred to as CLOUDS and NO CLOUDS, respectively. Model output consists of both 1D ( $t, z$ ) horizontally averaged 5-min statistics and 3D ( $t, x, y, z$ ) 5-min instantaneous data at the grid scale. The 1D data include contributions toward vertical fluxes from the subgrid-scale model; the 3D data do not. A comparison of 30-min averaged 1D and 3D data reveal no significant differences vertically across the lower troposphere (Figure S4 in Supporting Information S1). Consequently, both data sets are used throughout the following analysis.

### 2.3. Budget Analysis

To quantify the contributions of shallow convective clouds, clear air entrainment and the rainforest CO<sub>2</sub>-assimilation to the diurnal (daytime, between sunrise (6 LT) and sunset (18 LT)) and vertical (0–5 km) variability

of CO<sub>2</sub> above the Amazon, we study various decompositions of the atmospheric CO<sub>2</sub> budget with our LES simulations during the day. These analyses depart from the horizontally averaged vertically resolved CO<sub>2</sub> budget in the Boussinesq approximation (Equation 1) (Betts et al., 1990; Ouwersloot et al., 2013; Stull, 1988). Here, the terms represent the CO<sub>2</sub>-tendency (first term), mean horizontal (second and third term) and vertical advection (fourth term) of CO<sub>2</sub>, and the vertical divergence of the CO<sub>2</sub> flux (fifth term). These terms balance at all levels above the surface in the absence of external sources and sinks such as photosynthesis or respiration. The different terms within Equation 1 and the subsequent equations in this section are presented as 30-min averages directly determined from the 1D and 3D simulation outputs.

$$\frac{\partial \overline{\text{CO}_2}}{\partial t} + \bar{u} \frac{\partial \overline{\text{CO}_2}}{\partial x} + \bar{v} \frac{\partial \overline{\text{CO}_2}}{\partial y} + \bar{w} \frac{\partial \overline{\text{CO}_2}}{\partial z} + \frac{\partial}{\partial z} \overline{w' \text{CO}_2'} = 0 \quad (1)$$

Since horizontal and vertical advections (subsidence) are neglected in our LES, Equation 1 can be simplified to Equation 2:

$$\frac{\partial \overline{\text{CO}_2}}{\partial t} = -\frac{\partial}{\partial z} \overline{w' \text{CO}_2'} + R \quad (2)$$

Here  $R$  represents a residual term, which appears in the following as a result of errors made by (a) time-interpolating  $\overline{\text{CO}_2}$  when estimating its time derivative, and (b) ignoring the subgrid contributions to  $\overline{w' \text{CO}_2'}$  in our a-posteriori estimates of the budget.

### 2.3.1. Clear-to-Cloudy Boundary Layer CO<sub>2</sub> Budget

To answer the first research question, we study the vertically averaged clear-to-cloudy boundary layer CO<sub>2</sub> budget over time (Section 3.2). With clear-to-cloudy we refer to clear boundary layers that become cloudy due to the formation of shallow convective clouds throughout the day. The formulation of the budget follows from vertically integrating Equation 2 (ignoring  $R$  at present) from the canopy top, in DALES taken to be the lower surface ( $z_c = 0$  m), to the top of the well-mixed boundary layer ( $z_i$ ) and dividing by  $z_i$ . For simplicity, the subcloud layer height is taken synonymous with  $z_i$  when shallow convective clouds are present, reducing additional complexity involved in analyzing the dynamics of the transition from unsaturated to saturated air (Albright et al., 2022). For DALES,  $z_i$  is set equal to the height of the minimum buoyancy flux. Defining  $\langle \cdot \rangle$  as the vertical average, we obtain Equation 3 for the budget of boundary layer CO<sub>2</sub> in DALES,  $\langle \overline{\text{CO}_2} \rangle$ , where  $\overline{(w' \text{CO}_2')}_s$  is the flux of CO<sub>2</sub> at  $z_c$ , which is proportional to the net ecosystem exchange estimated by the A-g<sub>s</sub> model, and  $\overline{(w' \text{CO}_2')}_e$  representing the net flux at  $z_i$ , the interface between the clear-to-cloudy boundary layer and free troposphere.

$$\frac{\partial \langle \overline{\text{CO}_2} \rangle}{\partial t} = \frac{1}{z_i} \int_{z_c}^{z_i} \frac{\partial \overline{\text{CO}_2}}{\partial t} dz = \frac{\overline{(w' \text{CO}_2')}_s - \overline{(w' \text{CO}_2')}_e}{z_i} \quad (3)$$

to further disentangle the processes controlling  $\langle \overline{\text{CO}_2} \rangle$ , we expand the fluxes  $\overline{(w' \text{CO}_2')}_s$  and  $\overline{(w' \text{CO}_2')}_e$  into a form more suitable for our subsequent analysis and physical interpretation. First,  $\overline{(w' \text{CO}_2')}_s$  can be expanded based on the CO<sub>2</sub>-assimilation ( $A$ ) and respiration ( $Resp$ ) rates at canopy scale (Equation 4), following the A-g<sub>s</sub> model formulation from Ronda et al. (2001). The canopy-scale net CO<sub>2</sub>-assimilation ( $A_{n,c}$ ) is obtained by upscaling the net CO<sub>2</sub>-assimilation of individual leaves ( $A_n$ ) (Equation 5). Here  $A_g$  represents the gross CO<sub>2</sub>-assimilation rate,  $R_d$  the dark respiration,  $g_{l,c}$  the stomatal conductance at the leaf level,  $C_i$  the internal CO<sub>2</sub> mole fraction within the leaf, and  $C_s$  the CO<sub>2</sub> mole fraction at the leaf surface equal to the well-mixed atmospheric mole fraction ( $C_s = \langle \text{CO}_2 \rangle$ ). Upscaling  $g_{l,c}$  yields the stomatal conductance at the rainforest canopy level ( $g_{c,c}$ ). Substituting  $A_{n,c}$  of Equation 5 into Equation 4 and neglecting the aerodynamic resistance yields an approximation of  $\overline{(w' \text{CO}_2')}_s$  (Equation 6). Here,  $(\langle \text{CO}_2 \rangle - C_i)$  can be noted as  $\Delta C_{\text{vegetation}}$  since it denotes the difference in CO<sub>2</sub> between the atmosphere and the vegetation. Since  $Resp$  is approximately constant in time (Figure S1 in Supporting Information S1), the stomatal CO<sub>2</sub>-assimilation is the primary factor regulating the diurnal variability in  $\overline{(w' \text{CO}_2')}_s$ .



$$\overline{(w'CO_2')}_s = A_{n,c} + Resp \quad (4)$$

$$\overbrace{A_n = A_g - R_d = g_{lc}(C_s - C_i)}^{\text{leaf}} \rightarrow \overbrace{A_{nc} = g_{cc}(\langle CO_2 \rangle - C_i)}^{\text{canopy}} \quad (5)$$

$$\overline{(w'CO_2')}_s = A_{n,c} + Resp = g_{c,c}(\langle CO_2 \rangle - C_i) + Resp \approx g_{c,c} \Delta C_{vegetation} \quad (6)$$

Next, we expand  $\overline{(w'CO_2')}_e$ . Since we wish to quantify the contribution of dry and moist convective processes to  $\overline{(w'CO_2')}_e$ , we partition the flux by subsetting the DALES grid into the average contribution from clouds and the environment following the “top-hat” approximation commonly used in large scale numerical modeling (Siebesma & Cuijpers, 1995). More concretely, we partition  $\overline{(w'CO_2')}_e$  by selecting all vertical grid columns with a cloud present anywhere in the column (superscript *cl*, the in-cloud contribution), and all grid columns which do not have a cloud (superscript *env*, the environment contribution). Assessing the presence of a cloud in the vertical grid column is done using the 2D (*x, y*) cloud fraction of the DALES grid ( $a_x$ ) (Equation 7).

$$\overline{(w'CO_2')}_e = \overbrace{a_{cw} \overline{(w'CO_2')}_e^{cl}}^{\text{in-cloud}} + \overbrace{(1 - a_{cw}) \overline{(w'CO_2')}_e^{env}}^{\text{environment}} \quad (7)$$

It is important to note that the relative in-cloud and environment contributions to  $\overline{(w'CO_2')}_e$  highly depend on the selected cloud fraction. Selecting a different criterion defining the cloud fraction could yield different results. Throughout literature (e.g., Siebesma et al., 2003), two criteria are commonly presented for selecting cloud fractions within model grids: the “cloud-average” (liquid specific humidity ( $q_l$ ) > 0 g kg<sup>-1</sup>) and the “core-average” ( $q_l$  > 0 g kg<sup>-1</sup>, positive fluctuations in vertical velocity ( $w'$ ) > 0 m s<sup>-1</sup> and buoyancy ( $\theta'_v$ ) > 0 K compared to the domain average criterion. The cloud-average cloud fraction ( $a_c$ ) and core-average cloud fraction ( $a_{cc}$ ) are thereby related through approximately a factor 2 ( $a_c \approx 2a_{cc}$ ) (Sikma & Ouwersloot, 2015).

We anticipate that the ventilation of CO<sub>2</sub> by shallow convective clouds is predominantly driven by actively buoyant, upward-moving regions (updraughts) concentrated within the cloud cores. The dynamics of these buoyant cores can, in turn, generate additional small-scale turbulent motions along their edges, promoting mixing with surrounding neutral or negatively buoyant air through entrainment and detrainment processes (de Rooy et al., 2013; Savre, 2022). Consequently, larger volumes of air, and CO<sub>2</sub>, can be transported upward. To account for this process, we introduce a less strict criterion termed the “upward-moving-cloud-average” criterion ( $q_l$  > 0 g kg<sup>-1</sup> and  $w' > 0$  m s<sup>-1</sup>), with the upward-moving cloud-average fraction referred to as  $a_{c,w}$ . Adopting a different criterion primarily impacts how many grid columns are classified as cloud or environment, introducing a sensitivity in the overall contribution by in-cloud and environment grid cells. For completeness, we present the sensitivity of using the different cloud fractions to our main results in Figure S5 in Supporting Information S1.

For more quantitatively associating the partitioned fluxes in Equation 7 with in-cloud venting and clear air entrainment, we rewrite the equation using a well-mixed zero-order jump model formulation over an infinitesimal small layer similar to Equation 6 (Pino et al., 2006). We rewrite the terms as a product of a characteristic vertical velocity (entrainment), and the difference in CO<sub>2</sub> mole fraction between the well-mixed clear-to-cloudy boundary layer mole fraction ( $\langle CO_2 \rangle$ ) and the free tropospheric mole fraction at 100 m above  $z_i$  ( $C_{ft} \approx C_{z_i+100 \text{ m}}$ ) (similar to Albright et al., 2022):  $\Delta C_{atm} = C_{ft} - \langle CO_2 \rangle$ . We assume  $a_{c,w}$  to be negligibly small ( $1 - a_{c,w} \approx 1$ ) as the organized turbulence dominates the net flux (Siebesma & Cuijpers, 1995). Consequently, the expression for the clear air contribution is reduced to Equation 8, where we diagnose the clear air entrainment velocity  $w_e$  as the vertical velocity needed to satisfy the equation.

$$(1 - a_{c,w}) \overline{(w'CO_2')}_e^{env} = -w_e \Delta C_{atm}, \quad (8)$$

we write the in-cloud contribution as Equation 9, where we define the convective mass flux ( $M = a_{c,w} w_{cloud}$ ) (e.g., van Stratum et al., 2014), and again define  $w_{cloud}$  as the vertical velocity needed to satisfy Equation 9.

$$a_{c,w}(\overline{w'CO_2^{cl}})_e = a_{c,w}w_{cloud}\Delta C_{atm} = M\Delta C_{atm}, \quad (9)$$

taken together, Equations 8 and 9 yield a rewritten and extended version of Equation 7 which we use in Section 3.2 (Equation 10):

$$\overline{w'CO_2^I}_e = M\Delta C_{atm} - w_e\Delta C_{atm} \quad (10)$$

By substituting Equations 6 and 10 into Equation 3, we obtain Equation 11:

$$\frac{\partial \langle CO_2 \rangle}{\partial t} = \underbrace{\frac{Resp}{z_i}}_{(1)} - \underbrace{\frac{g_{c,c}}{z_i}}_{(2)} \Delta C_{vegetation} - \left( \underbrace{\frac{M}{z_i}}_{(3)} - \underbrace{\frac{w_e}{z_i}}_{(4)} \right) \Delta C_{atm} \quad (11)$$

This equation distinguishes between four main terms which influence the diurnal evolution of  $\langle CO_2 \rangle$ : respiration (a), rainforest  $CO_2$ -assimilation (b), in-cloud ventilation (c), and clear air entrainment (d). Since the respiration term (term 1) is kept constant throughout the diurnal cycle, we mainly examine the relative contribution of terms 2 to 4 as the simulation progresses.

### 2.3.2. Vertically Resolved $CO_2$ Budget and $CO_2$ - $H_2O$ Correlation

To study the role of the shallow convection throughout the lower tropical troposphere (Section 3.3), we extend the use of Equation 7 to evaluate the in-cloud and environment contributions to  $\overline{(w'CO_2^I)}$  at all model levels. Thereby, we still keep the 2D  $a_{c,w}$  as the projected cloud cover to define in-cloud and environment columns over the DALES grid. Inserting this decomposition into the vertically resolved  $CO_2$  budget Equation 2 then allows diagnosing their contributions to  $\overline{CO_2}(z)$ .

Finally, to assess whether shallow convective clouds and clear air entrainment play a role in organizing the transport of key variables in the rainforest-cloud system as represented in Figure 1 (Section 3.4) we partition the horizontally averaged covariance between  $H_2O$  (specific humidity,  $q$ ) and  $CO_2$  ( $\overline{(q'CO_2^I)}$ ) into in-cloud and environment contributions (Equation 12).

$$\overline{(q'CO_2^I)} = a_{c,w}(\overline{q'CO_2^{cl}}) + (1 - a_{c,w})(\overline{q'CO_2^{env}}) \quad (12)$$

Using Equation 12, we also calculate the partitioned horizontally averaged correlation factor ( $r$ ) by scaling  $\overline{(q'CO_2^I)}$  by the horizontal variation in  $q$  ( $\sigma_q$ ) and  $CO_2$  ( $\sigma_{CO_2}$ ) (Equation 13).

$$r = \frac{a_{c,w}(\overline{q'CO_2^{cl}}) + (1 - a_{c,w})(\overline{q'CO_2^{env}})}{\sigma_q \sigma_{CO_2}} \quad (13)$$

## 3. Results

### 3.1. Model Evaluation

We begin by evaluating the DALES simulation with the CloudRoots-Amazon22 observations, focusing on the most relevant measured quantities that affect the simulated energy, water, and carbon cycles above the rainforest: the surface energy balance,  $\overline{(w'CO_2^I)}_s$ , clear-to-cloudy boundary layer ( $z_i$ ), and cloud layer development, and the evolution of the vertical structure of specific humidity ( $q$ ), potential temperature ( $\theta$ ) and atmospheric  $CO_2$ .

### 3.1.1. Model Versus Observations: Surface Fluxes and Cloud Development

The simulated and observed diurnal evolution of the surface fluxes of sensible (H) and latent (LE) heat (Figure 2a) of the 6-day aggregate at 10 m above the canopy top ( $z = 50$  m) are similar, although DALES tends to slightly overestimate both (Root Mean Square Error (RMSE)—LE:  $66.76 \text{ W m}^{-2}$ , H:  $19.07 \text{ W m}^{-2}$ ), largely consistent with the net radiation available (Figure S1 in Supporting Information S1). Both observations and simulation suggest that the magnitude of LE is almost double the magnitude of H, characteristic for rainforest conditions with a high moisture content (e.g., von Randow et al., 2004).

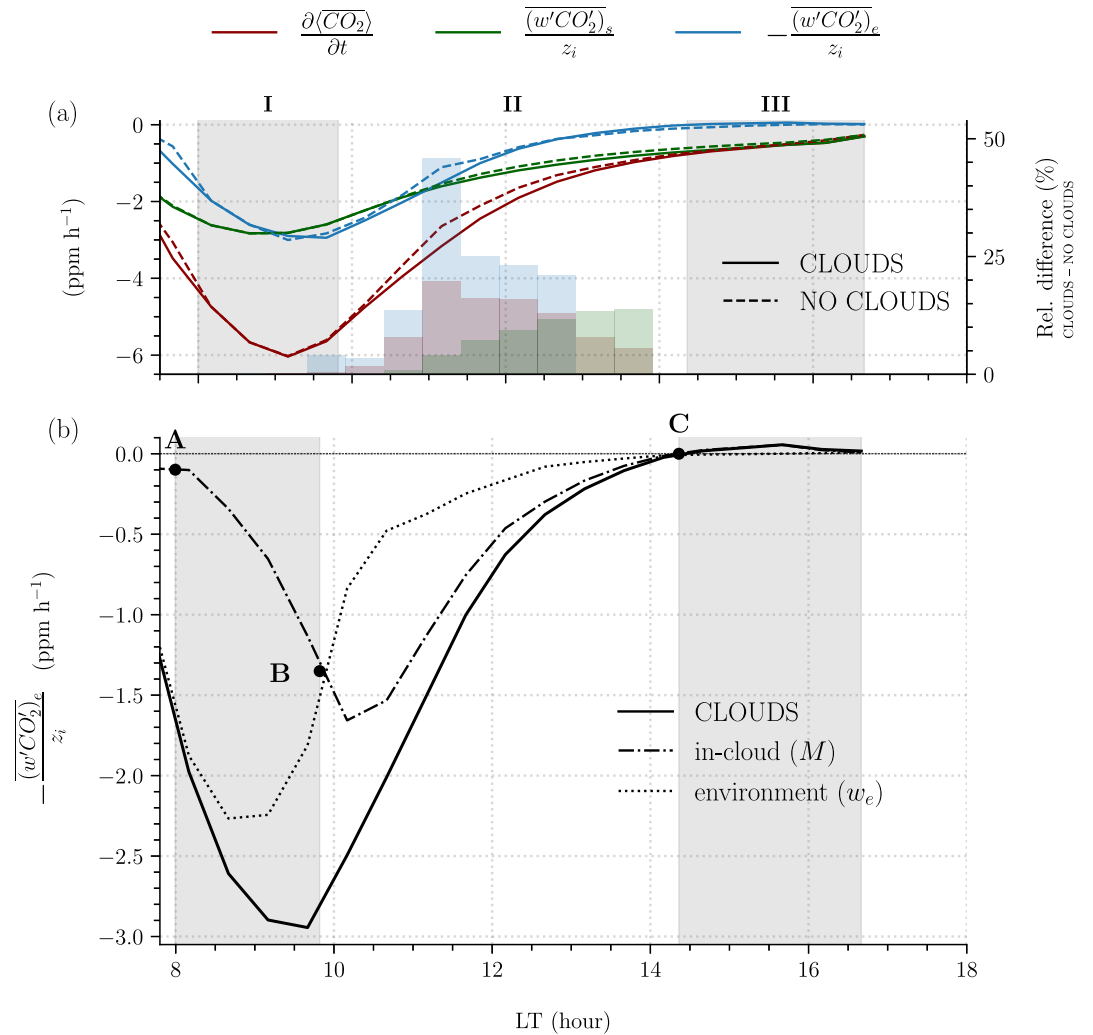
Similar to H and LE, DALES effectively simulates the diurnal cycle of  $(w'CO_2)_s$  (Figure 2b) (RMSE of  $5.45 \mu\text{mol m}^{-2} \text{ s}^{-1}$ ), with a clear photosynthesis signal but with some shifts in the minimum values throughout the day. Throughout the day, the rainforest has an uptake of  $10\text{--}20 \mu\text{mol m}^{-2} \text{ s}^{-1}$  ( $\approx 18\text{--}36 \text{ ppm h}^{-1}$  uptake throughout the canopy ( $z = 50$  m) in standard atmospheric conditions  $T_0 = 25^\circ\text{C}$  and  $p_0 = 1,013.25 \text{ hPa}$ ). A notable difference is the observed near-zero flux during the morning transition (6–8 LT), which is not represented within DALES. González-Armas et al. (2025), who investigated the in-canopy fluxes observed during CloudRoots-Amazon22, suggests that there is a complex flux divergence profile within the canopy that our single-layer, bulk canopy surface model does not appropriately account for. This flux divergence results from a combination of the flush-up of stored  $CO_2$  within the canopy, the initiation of photosynthesis at the canopy top, and ongoing respiration at the canopy bottom, as the subcanopy remains decoupled from the upper canopy and atmosphere (Dupont et al., 2024). Combined, these features offset the measured flux, resulting in a net zero flux as observed during the morning transition. Missing these features in the bulk canopy representation may explain why DALES cannot produce a near zero-flux during the morning transition. Fortunately, after 8 LT, once the model has spun-up and the morning transition has finished, the simulated  $(w'CO_2)_s$  becomes more realistic. Hence, we can focus our analysis after this time step.

The development of  $z_i$ , clouds and cloud layer (Figure 2c) are also accurately represented by DALES (RMSE  $z_i$  of 75.63 m) and in accordance to previous research (Giangrande et al., 2020; Henkes et al., 2021; Vilà-Guerau de Arellano et al., 2020). After a stable early morning,  $z_i$  grows rapidly by  $150\text{--}200 \text{ m h}^{-1}$ . After 10 LT, active shallow convective clouds form as  $a_c$ ,  $a_{c,w}$  and  $a_{cc}$  increase (Figure 2d), peaking at 0.30, 0.23, and 0.16 at 14 LT, respectively. The formed clouds have a base around 1.5 km, with a horizontal and vertical extent of respectively 500 m and 0.5–2 km, in both model and observations. After 14 LT, clouds decrease in activity and gradually dissipate toward a nearly cloud-free dusk. The evolution of the cloud layer top, and consequently the vertical extent to which the clouds can influence the vertical distribution of  $CO_2$ , follows the expected linear time-scaling with surface buoyancy flux (Stevens, 2007).

### 3.1.2. Model Versus Observations: Vertical Structure of the Lower Troposphere

In the vertical, DALES successfully captures  $\theta$ ,  $q$  and  $CO_2$  in the lower tropical troposphere at 12 and 13 LT (Figures 3a–3c, vertical profiles at 6, 9, 12, 15, and 18 LT are included in Figure S2 in Supporting Information S1). For  $\theta$ , a well-mixed layer is clearly present in the lowest 1.5 km. The model and observations align well, with a RMSE of 0.45 K. Similar to  $\theta$ ,  $q$  is observed to be relatively well-mixed below 1,200 m. However, the LES overestimates this well-mixedness somewhat, yielding a RMSE of  $1.12 \text{ g kg}^{-1}$ . Similar to  $\theta$  and  $q$ , DALES successfully captures the vertical structure of  $CO_2$  (RMSE: 0.94 ppm), with a small inversion around 1,200 m at 13 LT.

The 6-day aggregated tower observations of the  $CO_2$  mole fractions at 24, 38, 53, 79, and 321 m reflect the aforementioned (Section 2.2) flush-up signal after sunrise (Culf et al., 1997; Dupont et al., 2024; González-Armas et al., 2025) (Figure 3d). These observations show that the relatively high  $CO_2$  mole fractions near the surface (24 m) decrease rapidly, whereas mole fractions higher up the tower (38–79 m) increase at a similar rate. At 53 m, which is  $\sim 13$  m above the canopy,  $CO_2$  mole fractions remain relatively constant before 10 LT, reflecting a well-mixed value composed of the values from in-canopy air with the residual layer above. The  $CO_2$  mole fraction at 321 m, the furthest observation away from the canopy and thus not directly influenced by the rainforest processes, approaches this mixed-layer value just before 10 LT, two to three hours after the lowest 80 m of the atmosphere. The moment at which the 321 m level reaches this well-mixed value coincides with the effective onset of shallow convective clouds and a marked increase in photosynthesis (Figure 2b), followed by a  $\sim 20 \text{ ppm}$  reduction toward the mid-afternoon.

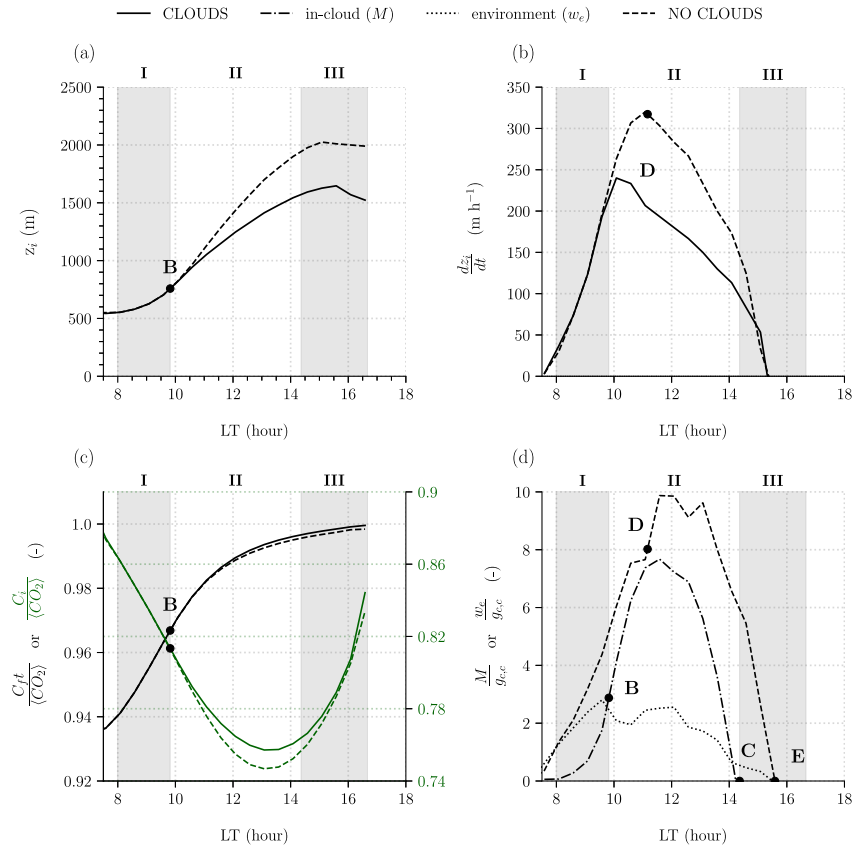


**Figure 4.** Atmospheric CO<sub>2</sub> budget in the clear-to-cloudy boundary layer (Equation 3) for (a) the CLOUDS (solid line) and NO CLOUDS (dashed line) simulation between 6 and 18 LT, along with the relative difference (%), represented as bars) between the two simulations ( $(\text{CLOUDS} - \text{NO CLOUDS})/\text{NO CLOUDS}$ ), relative to the NO CLOUDS simulation. Additionally, (b) shows the in-cloud (dashed dotted line) and environment (dotted line) contributions (Equation 7) to  $\langle w' \text{CO}_2' \rangle_e$  of the CLOUDS simulation. The numbers I through III and shading denote the three regimes that can be distinguished from left to right: entrainment-diluting regime (I), cloud-ventilation-and-entrainment regime (II) and the CO<sub>2</sub>-assimilation regime (III). Markers A through E denote key time steps aiding in the description of the diurnal cycle of both simulations.

Since the 321 m observations represent the well-mixed layer value following the morning flush-up, we can use this height to assess whether DALES, which we constrained manually to match near-canopy observations (53 m) (Section 2.2), accurately simulates these mole fractions over time. Figure 3d illustrates that DALES effectively captures these mole fractions of CO<sub>2</sub> at 321 m from 8 LT onward, with a closer fit after 10 LT (RMSE CO<sub>2</sub> after 8 LT: 4.723 ppm, after 10 LT: 1.995 ppm), allowing us to assess the CO<sub>2</sub> budget from these times onward.

### 3.2. Clear-to-Cloudy Boundary Layer CO<sub>2</sub> Budget

To quantify the contributions from in-cloud venting and clear air entrainment alongside rainforest CO<sub>2</sub>-assimilation on the clear-to-cloudy boundary layer CO<sub>2</sub> mole fractions, we present the total tendency for the CLOUDS and NO CLOUDS simulations, and these contributions to the total tendency of the CLOUDS simulation, throughout the day (Equation 7) in, respectively, Figures 4a and 4b. Moreover, we present the relative difference (%) between the CLOUDS and NO CLOUDS simulation ( $(\text{CLOUDS} - \text{NO CLOUDS})/\text{NO CLOUDS}$ ), relative to the NO CLOUDS simulation, in Figure 4a. For our analysis, we break down each flux contribution following Equation 11 in



**Figure 5.** Different terms composing the atmospheric CO<sub>2</sub> budget in the clear-to-cloudy layer from Equation 11: (a) boundary layer height ( $z_i$ ), (b) growth rate of  $z_i$ , (c) ratios between the internal mole fraction ( $C_i$ ) or free tropospheric mole fraction ( $C_f$ ) and the clear-to-cloudy boundary layer mole fraction ( $\langle \text{CO}_2 \rangle$ ) of CO<sub>2</sub>, and (d) ratios of the characteristic velocities (in-cloud (dashed dotted line) contribution represented by the mass flux ( $M$ ), the environment (dotted line) contribution and NO CLOUDS (dashed line) simulations represented by the entrainment velocity ( $w_e$ ), against the stomatal conductance at the rainforest canopy level ( $g_{c,c}$ ). The numbers I through III and shading denote the three regimes that can be distinguished from left to right: entrainment-diluting regime (I), cloud-ventilation-and-entrainment regime (II) and the CO<sub>2</sub>-assimilation regime (III). Markers A through E denote key time steps aiding in the description of the diurnal cycle of both simulations.

Figure 5. More specifically, we present the development and growth rate of  $z_i$  (Figures 5a and 5b), the ratio between the clear-to-cloudy boundary layer and the cloud layer ( $\frac{C_i}{\langle \text{CO}_2 \rangle}$ ) or internal CO<sub>2</sub> mole fractions ( $\frac{C_f}{\langle \text{CO}_2 \rangle}$ ) resembling  $\Delta C_{\text{atm}}$  and  $\Delta C_{\text{vegetation}}$  (Figure 5c), and the ratios of the characteristic velocities ( $M$  and  $w_e$  against  $g_{c,c}$ ) (Figure 5d) in time. The markers A through E in both figures denote key time steps aiding in the description of the diurnal cycle of both simulations. In the following, we first concentrate on analyzing the CLOUDS simulation (solid lines) (Section 3.2.1), before assessing the differences with the budget in the NO CLOUDS simulation (Section 3.2.2).

### 3.2.1. With Clouds: Joint Regime of Cloud Ventilation and Entrainment

Throughout the day, CO<sub>2</sub> is ventilated and mixed at a rate ( $\frac{\partial \langle \text{CO}_2 \rangle}{\partial t}$ ) of 2–6 ppm h<sup>−1</sup>, diluting the clear-to-cloudy boundary layer in CO<sub>2</sub> (Figure 4a). Although both  $\overline{(w' \text{CO}_2')_s}$  and  $\overline{(w' \text{CO}_2')_e}$  contribute to this dilution, the latter outweighs the former before 14:21 LT, establishing an atmospheric-dilution regime and CO<sub>2</sub>-assimilation regime, respectively, in accordance to previous research (Davis et al., 1997; Faassen et al., 2025; Huang et al., 2011; Verzijlbergh et al., 2009; Vilà-Guerau de Arellano et al., 2004). However, partitioning  $\overline{(w' \text{CO}_2')_e}$  between in-cloud (dash-dotted lines) and environment (dotted lines) contributions according to Equation 7 reveal that with the presence of clouds,  $\overline{(w' \text{CO}_2')_e}$  shifts from being dominated by clear air entrainment before 9:49 LT,



to having a joint contribution between in-cloud venting and clear air entrainment under increasing cloud fraction. As a result, we can separate the atmospheric-dilution into two separate regimes: an “entrainment-diluting regime” before 9:49 LT, and the “cloud-ventilation-and-entrainment regime” between 9:49 and 14:21 LT (previously hinted to by Faassen et al. (2025) and Verzijlbergh et al. (2009)).

The onset of the entrainment-diluting regime occurs after model initialization, flush-up and start of dry convection at 8 LT (marker A). At this time, a growing entrainment velocity  $w_e$  (Figure 5d) rapidly grows the boundary layer (markers A to B). The combined effect on the  $\langle \text{CO}_2 \rangle$  budget (term 4 in Equation 11) reduces  $\langle \text{CO}_2 \rangle$  at a rate of  $-3 \text{ ppm h}^{-1}$ , contributing roughly half the  $\langle \text{CO}_2 \rangle$  sink. The effectiveness of this entrainment-dilution is limited both by the growing  $z_i$ , and by the dilution itself, which brings  $\langle \text{CO}_2 \rangle$  toward  $C_{fi}$ . As a result, the ratio between these mole fractions ( $\frac{C_{fi}}{\langle \text{CO}_2 \rangle}$ , Figure 5c) approaches a value of one (0.985). Approximately 2 hours later (09:49 LT), following the formation of more actively buoyant shallow convective clouds (Figure 2d), the cloud-ventilation-and-entrainment regime commences (marker B). At this time, the contribution of in-cloud ventilation (term 3 in Equation 11) increases rapidly, whereas that of  $w_e$  (term 4) gradually decreases (Figure 5d).

At 14:21 LT, approximately 4 hr and 30 min after the onset of moist convection,  $(w' \text{CO}_2)_e$  ceases as the lower troposphere reaches well-mixed conditions ( $\frac{C_{fi}}{\langle \text{CO}_2 \rangle} > 0.995$ , marker C). At this moment, the third and last regime, the  $\text{CO}_2$ -assimilation regime, commences as rainforest  $\text{CO}_2$ -assimilation (term 2 in Equation 11) becomes the sole controlling factor of the  $\text{CO}_2$  budget of the boundary layer. Apart from the  $\text{CO}_2$ -assimilation regime, photosynthesis is an active sink of  $\text{CO}_2$  throughout the entire diurnal cycle. Assimilation of  $\text{CO}_2$  has a similar magnitude as  $(w' \text{CO}_2)_e$ , despite the fact that the canopy-scale conductance is significantly smaller than the mass flux and entrainment velocities ( $M$  and  $w_e$  being 2 to 8 times larger, as shown in Figure 5d). The effectiveness of the  $\text{CO}_2$ -assimilation by the vegetation is primarily governed by the fact that the gradient in  $\text{CO}_2$  for the vegetation ( $\Delta C_{\text{vegetation}}$ ) is far greater than that between the boundary layer and free troposphere ( $\Delta C_{\text{atm}}$ ). The vegetation actively uses  $C_i$  for photosynthesis, reducing  $\frac{C_i}{\langle \text{CO}_2 \rangle}$  until midday while maintaining  $\text{CO}_2$  uptake as the boundary layer  $\text{CO}_2$  mole fraction decreases. Near 13 LT, the internal mole fraction of  $\text{CO}_2$  for the vegetation is lowered to  $\sim 25\%$  the well-mixed boundary layer value. This behavior is in accordance with observations and detailed analysis presented by González-Armas et al. (2025).

Overall, throughout the entire day, the clear-to-cloudy boundary layer  $\text{CO}_2$  budget is controlled by an interplay between in-cloud ventilation ( $\sim 23\%$ ), clear air entrainment ( $\sim 21\%$ ), next to continued rainforest  $\text{CO}_2$ -assimilation ( $\sim 56\%$ ), with their relative importance varying per diurnal regime.

### 3.2.2. Without Clouds: Entrainment Becomes the Governing Process

When latent heating is prevented from influencing buoyancy (NO CLOUDS, dashed lines), only two regimes remain: the entrainment-diluting and the  $\text{CO}_2$ -assimilation regime. However, the overall effect on  $\frac{d\langle \text{CO}_2 \rangle}{dt}$  appears visually minimal (Figure 4a). An examination of the different terms comprising the atmospheric  $\text{CO}_2$  budget in the clear-to-cloudy layer, as expressed in Equation 11, shows that the difference between the two simulations is primarily governed by subtle compensating effects between  $M$ ,  $z_i$ , and  $w_e$ , which we shall discuss in more detail below.

Within the NO CLOUDS simulation,  $M$  is forced to zero as clouds are considered passive. As a consequence, cloud ventilation (term 3 in Equation 11) is lost. If the missing cloud ventilation is treated individually, no ventilation by mass fluxes should yield a relative accumulation of  $\langle \text{CO}_2 \rangle$  compared to the CLOUDS simulation. However, a striking observation is that without clouds, increasing  $w_e$  (term 4 in Equation 11) compensates for the missing cloud ventilation.

Next to the missing  $M$ ,  $z_i$  becomes up to 500 m larger for the NO CLOUDS simulation compared to the CLOUDS simulation, comparable to Vilà-Guerau de Arellano et al. (2005). Shallow convective clouds act as valves which evacuate mass from the subcloud layer (e.g., Nicholls & Lemone, 1980; van Stratum et al., 2014; Verzijlbergh et al., 2009). As a result of this mass removal,  $z_i$  grows about  $150 \text{ m h}^{-1}$  slower in the CLOUDS simulation compared to the NO CLOUDS simulation once active buoyant clouds form around 10 LT (marker B). The effects of the larger  $z_i$  is twofold. On the one hand, the increasing  $z_i$  causes  $\langle \text{CO}_2 \rangle$  to be diluted across a larger volume, reducing the overall  $\text{CO}_2$  mole fraction. On the other hand, the increasing volume reduces the entrainment-

dilution in the NO CLOUDS simulation as every entrained  $\text{CO}_2$ -dilute parcel must be mixed into a deeper layer to lower  $\langle \text{CO}_2 \rangle$ . If the effects of the growing  $z_i$  would be treated individually, an increasing  $z_i$  should lead to a lower  $\langle \text{CO}_2 \rangle$  when clouds are passive. However, similar to  $M$ , the effects are strikingly compensated by the increase in  $w_e$  in the NO CLOUDS simulation.

A possible explanation of the increased  $w_e$  can be found in the radiation and energy balance, as the rainforest canopy top receives  $55 \text{ W m}^{-2}$  more net radiation under clear sky conditions. Consequently, the sensible heat flux is increased by  $32 \text{ W m}^{-2}$  and the development of large, strong turbulent structures is promoted, increasing  $w_e$ . Moreover, since more grid columns are considered as “clear”, the relative contribution from the environment and thus  $w_e$  in Equation 11 increases. It is worth noting that  $w_e$  peaks around midday after a fast increase at 11 LT (marker D in Figure 5c).

The increase in  $w_e$  is the primary process responsible for yielding an approximately similar NO CLOUDS and CLOUDS simulation throughout the day. However, quantifying the relative difference between the CLOUDS and NO CLOUDS simulations, relative to the NO CLOUDS simulation, reveals a substantial discrepancy between 10 and 14 LT. During this period, the total tendency in the CLOUDS simulation is approximately 20%–25% greater (Figure 4a). The largest contribution to this discrepancy arises from  $(w' \text{CO}_2')_e$ , which is more than 25% greater in the CLOUDS simulation following the onset of shallow convective clouds. Although the increasing  $w_e$  and  $z_i$  in the NO CLOUDS simulation compensate to some extent for the absence of  $M$ , the relative difference highlights that the net removal by  $(w' \text{CO}_2')_e$  in the NO CLOUDS simulation is 20%–25% less effective than when clouds are present. As a result, the NO CLOUDS simulation reaches well-mixed conditions approximately an hour later, at 15:35 LT (marker E). Additionally, the rainforest vegetation must lower its  $C_i$  to adjust to the slightly higher  $\langle \text{CO}_2 \rangle$  between 10 and 14 LT (Figure 5c).

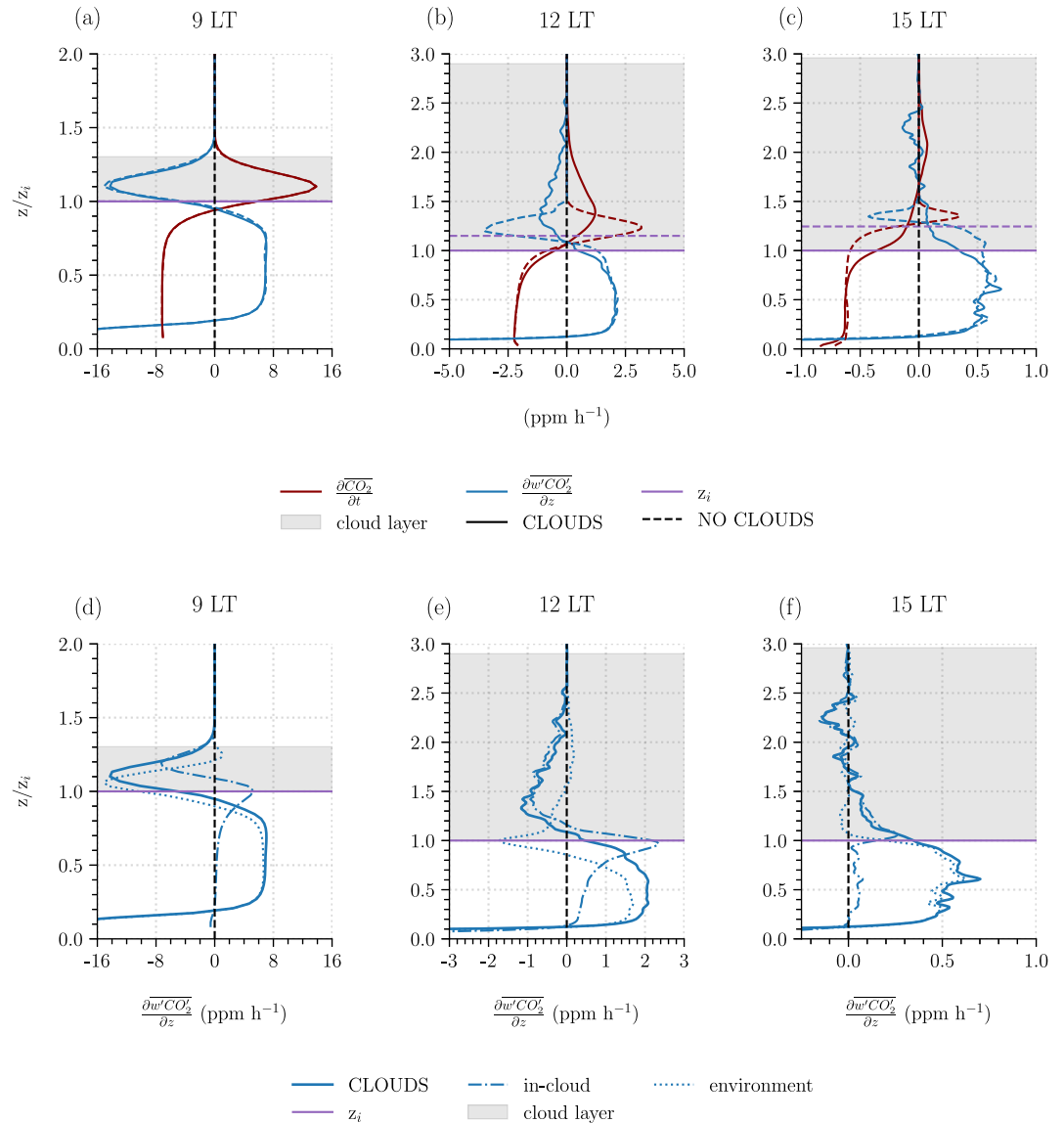
### 3.3. Vertically Resolved $\text{CO}_2$ Budget: Shallow Convective Clouds Raise Vertical $\text{CO}_2$ Levels Up to Double $z_i$

The CLOUDS and NO CLOUDS simulations differ more substantially above  $z_i$  during the cloud-ventilation-and-entrainment regime (Figure 6). Before this time, at 9 LT (Figure 6a), the vertically resolved  $\text{CO}_2$  budget follows a clear entrainment profile: dry convective thermals transport  $\text{CO}_2$  from the canopy and boundary layer (net divergence of  $8 \text{ ppm h}^{-1}$ ) to a shallow layer around  $z_i$  (between  $z_i$  to  $1.3z_i$ ), where  $\text{CO}_2$  is deposited (net convergence of  $16 \text{ ppm h}^{-1}$ ). At the same time,  $\text{CO}_2$ -assimilation by the rainforest actively removes  $\text{CO}_2$ , explaining the negative  $\text{CO}_2$  tendency down to the lowest model level.

At midday (Figure 6b), the entrainment signal as observed at 9 LT, though in smaller magnitude ( $2.5 \text{ ppm h}^{-1}$ ), remains for the NO CLOUDS simulation. For the CLOUDS simulation, the convergence signal of entrainment is spread up to  $2.5z_i$ , following the linear time-scaling with surface buoyancy flux (Stevens, 2007). The signal has a characteristic bell-shaped distribution (Ouwensloot et al., 2013; Verzijlbergh et al., 2009; Vilà-Guerau de Arellano et al., 2005), with the largest convergence of  $\text{CO}_2$  with a rate up to  $2 \text{ ppm h}^{-1}$  at the bottom of the cloud layer, and decreasing values with height. Flux partitioning into actively buoyant in-cloud and environment grid columns according to Equation 7 (Figure 6e), reveal that this bell-shaped convergence signal can be primarily attributed to cloudy grid columns.

The vertical distribution of  $\text{CO}_2$  results from an interplay between entrainment from the boundary layer and in-cloud ventilation. Starting at the surface, dry turbulence in clear air (environment grid columns) transports  $\text{CO}_2$  upward (divergence of  $2 \text{ ppm h}^{-1}$ ). At  $z_i$ , the transported  $\text{CO}_2$  is deposited (convergence of  $2 \text{ ppm h}^{-1}$ ), but this is balanced by a net divergence of slightly more than  $2 \text{ ppm h}^{-1}$  in the in-cloud grid columns. That is, during the cloud-ventilation-and-entrainment regime, primarily dry convective thermals bring  $\text{CO}_2$  from the subcloud layer to  $z_i$ , whereas at  $z_i$ , moist convective thermals pick up this  $\text{CO}_2$  and transport it to higher vertical levels within the cloud layer, increasing  $\text{CO}_2$  mole fractions up to  $2.5z_i$ .

Toward the end of the afternoon during the  $\text{CO}_2$ -assimilation regime (Figures 6c and 6f), the signals from both simulations weaken as the lower troposphere reaches a well-mixed state. As a result, there is minimal to no net vertical displacement of  $\text{CO}_2$  within the lower troposphere. Without clear convergence at  $z_i$ , the remaining  $\text{CO}_2$  divergence within the boundary layer can be primarily attributed to continued  $\text{CO}_2$ -assimilation. However, since

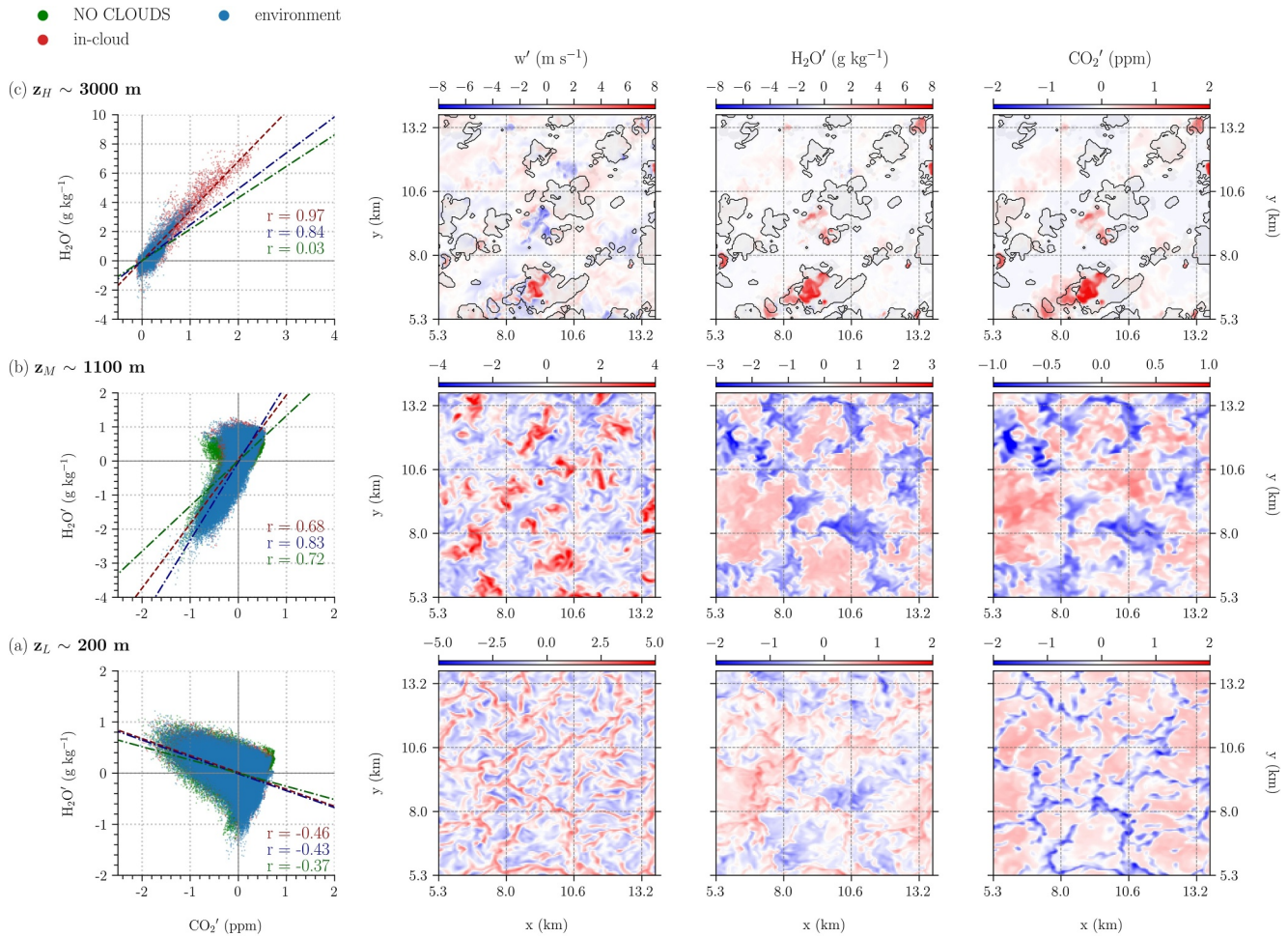


**Figure 6.** Vertically resolved budget of atmospheric CO<sub>2</sub> for the NO CLOUDS and CLOUDS (solid line) simulations (a–c), split into in-cloud (dashed dotted line) and environment (dotted line) contributions (d, e), at three time steps: (a, b) 9 LT, (b, e) 12 LT and (c, f) 15 LT.

the cloud fraction is still larger than zero (Figure 2d), a weak (0.2 ppm h<sup>−1</sup>) cloud ventilation signal remains visible, transporting CO<sub>2</sub> upward and primarily mixing the atmosphere >2.5z<sub>i</sub>.

### 3.4. Shallow Convective Clouds Organize the Vertical Correlation Pattern Between CO<sub>2</sub> and H<sub>2</sub>O

The preceding analysis of the contributions to the boundary layer (Section 3.2) and the vertically resolved CO<sub>2</sub> budget (Section 3.3) highlight the important role of shallow convective clouds in shaping these budgets. The associated mixing and vertical transport processes point to a pronounced coupling between the rainforest surface and the upper atmosphere (extending up to 2.5z<sub>i</sub>), regulated by these clouds. The presence of such coupling reflects the observed transition from a negative to positive CO<sub>2</sub>-H<sub>2</sub>O correlation with height as shown in Figure 1. To determine the extent to which shallow convective clouds shape this coupling, we proceed to analyze the CO<sub>2</sub>-H<sub>2</sub>O correlation signal and the horizontally averaged correlation coefficient (*r*) between fluctuations of H<sub>2</sub>O and CO<sub>2</sub> at z<sub>L</sub>, z<sub>M</sub>, and z<sub>H</sub> at 13 LT in DALES. We compare the NO CLOUDS simulation with the in-cloud and environmental contributions of the CLOUDS simulation (together they represent the CLOUDS signal), alongside



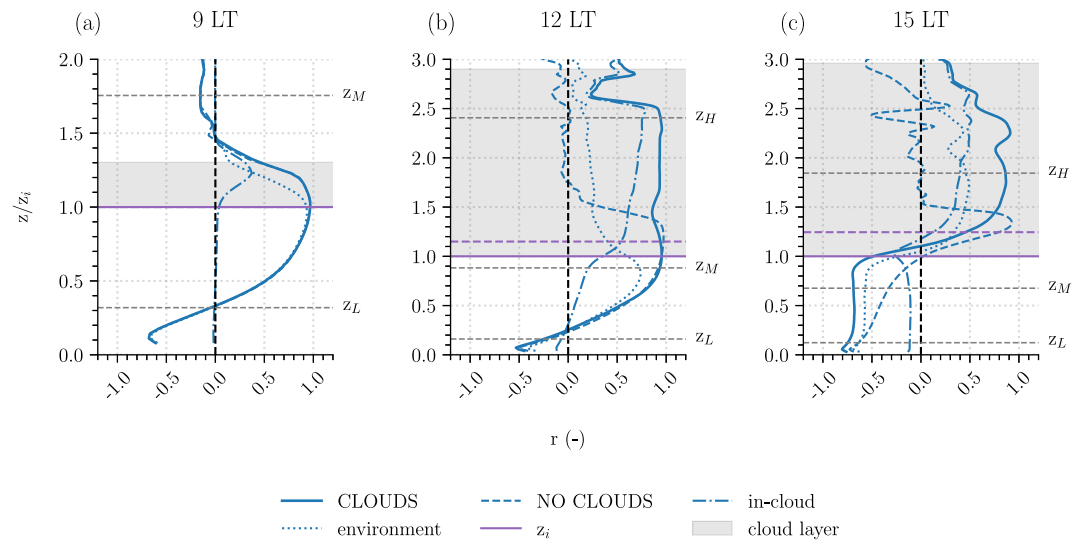
**Figure 7.** Horizontal instantaneous  $\text{CO}_2$ - $\text{H}_2\text{O}$  correlation signal and the horizontally averaged correlation coefficient ( $r$ ) in DALES at 13 LT between fluctuations of  $\text{H}_2\text{O}$  and  $\text{CO}_2$  for the NO CLOUDS and in-cloud and environmental contributions of the CLOUDS simulation (together they represent the CLOUDS signal) (left column), alongside the instantaneous horizontal ( $x, y$ ) cross section of fluctuations ( $\varphi'$ ) in  $\text{H}_2\text{O}$ ,  $\text{CO}_2$  and vertical velocity ( $w$ ) around spatial domain-averages ( $\bar{\varphi}$ ) in the CLOUDS simulation with the upward-moving cloud average boundaries colored in black contours ( $a_{c,w}$ ) (right columns). Represented at  $z_L$  ((a), roughness sublayer  $\sim 200$  m),  $z_M$  ((b), subcloud layer  $\sim 1,100$  m),  $z_H$  ((c), cloud layer at  $\sim 3,000$  m) similar to Figure 1. Simulated mean cloud-average cloud fraction ( $a_c$ ), upward-moving cloud-average cloud fraction ( $a_{c,w}$ ) and core-average cloud fraction ( $a_{cc}$ ) are approximately: 0.28, 0.21, and 0.14, respectively.

the instantaneous horizontal ( $x, y$ ) cross section of the CLOUDS simulation with the upward-moving cloud average boundaries ( $a_{c,w}$ ) in Figure 7. The instantaneous cross section represents a spatial representation of  $\text{H}_2\text{O}'$  and  $\text{CO}_2'$ , similar to the flight tracks as shown Figure 1. Benefiting from the LES simulation, we can also assess the fluctuations in  $w$  to get the full picture. In the following, we first examine the CLOUDS simulation, before comparing it with the NO CLOUDS simulation. We thereby assess the vertical aspects prior to considering the temporal aspects.

The correlation figures and horizontally averaged  $r$ -values of DALES at 13 LT reflect the negative to positively correlated pattern with height, similar to Figure 1. However, in contrast to the observations, the positive correlation is already clearly defined at  $z_M$ , rather than solely at  $z_H$ . This difference may be attributed to minor discrepancies between the DALES simulation and the observations, as discussed in Section 3.1. Additionally, the absence of surface heterogeneity from the Amazon rainforest within DALES, which is captured by the aircraft observations, may also contribute to this difference. As a result, we focus on comparing the processes occurring at the roughness sublayer ( $z_L$ ) and cloud layer ( $z_H$ ).

At  $z_L$  (Figure 7a), the strong negative correlation ( $r \approx -0.5$ ) can be attributed to the mixing of  $\text{CO}_2$ -diluted air from photosynthesis and moisture-rich air from evapotranspiration by the rainforest in upward moving thermals, which is in accordance to previous research (Dupont et al., 2024; Vilà-Guerau de Arellano et al., 2004) and





**Figure 8.** Vertical profiles of the horizontally averaged correlation coefficient ( $r$ ) between  $\text{CO}_2$  and  $\text{H}_2\text{O}$  for the CLOUDS (solid line), split into in-cloud (dashed dotted line) and environment (dotted line) contributions, at three time steps: (a) 9 LT, (b) 12 LT and (c) 15 LT. Vertical lines highlight the heights of  $z_L$  (roughness sublayer at  $\sim 200$  m),  $z_M$  (subcloud layer at  $\sim 1,100$  m) and  $z_H$  (cloud layer at  $\sim 3,000$  m) similar to Figure 1.

observations from Figure 1. At  $z_H$  (Figure 7c), the strong positive correlation ( $r \approx 1$ ) can be attributed to an interplay of in-cloud ventilation and the environment (Jonker et al., 2008; Ouwersloot et al., 2013; Verzijlbergh et al., 2009). This interplay is twofold. On the one hand, convective thermals originating at the surface, carrying bulk boundary-layer moisture and  $\text{CO}_2$ , possess an excess of both  $\text{CO}_2$  and moisture compared to decreasing ambient values with height. This excess in  $\text{CO}_2$  and  $\text{H}_2\text{O}$  is jointly transported upward, reflecting the positive correlation. In dry thermals, both  $\text{CO}_2$  and  $\text{H}_2\text{O}$  are deposited close to  $z_i$ , whereas in cloudy thermals  $\text{CO}_2$  and  $\text{H}_2\text{O}$  are transported up to  $2.5z_i$ . Figure 7c shows that clouds actively organize this vertical transport, with  $r$  increasing by 0.13 in the in-cloud columns. Conversely, in the environment, compensating descending motions from the free troposphere, characterized by reduced  $\text{CO}_2$  and moisture from the free troposphere, entrain into the boundary layer. Consequently, these motions also give rise to a positive correlation.

The spatial representation of  $w'$ ,  $\text{H}_2\text{O}'$  and  $\text{CO}_2'$  reveal that the vertical exchange of  $\text{CO}_2$  at  $z_H$  is confined to cloud scales of  $<2$  km. Within these dynamic clouds, strong updraughts ( $w'$  up to  $9 \text{ m s}^{-1}$ ) bring air with bulk-boundary layer moisture and  $\text{CO}_2$  mole fractions upward ( $\phi'$  up to  $3\text{--}6 \text{ ppm}$  or  $\text{g kg}^{-1}$ ). Depending on the core updraught strength,  $\text{CO}_2$  and  $\text{H}_2\text{O}$  are displaced laterally at the cloud base, halfway the cloud layer or at the cloud top. These cloud scales typically fall within the “gray zone” at which convection in numerical weather prediction and climatological models can be represented (Honnert et al., 2020; Pedruzo-Bagazgoitia et al., 2019; Tomassini et al., 2023), making it a challenge to accurately account for these processes in estimates of the diurnal  $\text{CO}_2$  budget.

Without clouds, the turbulent thermals at  $z_L$  grow and decrease in magnitude (Dupont et al., 2024; Hooff et al., 2019), eventually vanishing above the raised  $z_i$ , resulting in  $r = 0$  at  $z_H$ . These thermals detrain their  $\text{CO}_2$  and moisture in a layer between  $z_i$  and  $1.3z_i$  and cannot transport  $\text{CO}_2'$  and  $\text{H}_2\text{O}'$  to  $z_H$ , as discussed in Section 3.3. These findings confirm that shallow convective clouds are responsible for extending the positive  $\text{CO}_2\text{--H}_2\text{O}$  correlation above  $z_i$ , as the correlated patterns remain consistent throughout the boundary layer for the CLOUDS and NO CLOUDS simulation at  $z_L$  and  $z_M$  (Figures 7b and 7c).

Over the course of the day, from 9 to 15 LT (Figure 8), the  $\text{CO}_2\text{--H}_2\text{O}$  correlation pattern with height remains a distinct structure. At both 9 and 12 LT, the positive correlation in the NO CLOUDS simulation extends up to the top of the entrainment zone (up to  $1.5z_i$ ), whereas in the CLOUDS simulation, the positive  $r$  extends up to  $2.5z_i$  and is primarily attributed by cloudy grid columns. By 15 LT, the correlation is primarily organized by fluctuations in  $\text{H}_2\text{O}$  since the boundary layer becomes well-mixed in  $\text{CO}_2$ . For  $\text{H}_2\text{O}$ , a strong and progressively more pronounced inversion develops over the course of the day, shaped by the relatively dry free troposphere above and the



replenishment of moisture through evapotranspiration over the rainforest. Consequently, entrainment from the boundary layer and in-cloud ventilation primarily transports moist air upward and drier air downwards. It is worth noting that this transport is sufficiently strong that a positive correlation above  $z_i$  can still be observed.

#### 4. Conclusions

In this study, we investigated the turbulent exchange of  $\text{CO}_2$  in the lower troposphere above the Amazon rainforest across clear boundary layers that become cloudy due to the formation of shallow convective clouds throughout the day. Specifically, we examined the interactions between clear air entrainment (dry convection), in-cloud ventilation (moist convection), next to rainforest  $\text{CO}_2$ -assimilation. Furthermore, we examined whether shallow convective clouds organize the vertical  $\text{CO}_2$  exchange within the lower tropical troposphere, as illustrated in Figure 1. We constructed and validated a representative DALES simulation of a shallow convective day above the Amazon rainforest using the CloudRoots-Amazon22 campaign observations. Two simulations were run: one without and one with clouds, where we partitioned the latter into its respective in-cloud and environmental contributions to the clear-to-cloudy boundary layer and vertically resolved  $\text{CO}_2$  budget. Our results demonstrated that the interplay between clear air entrainment, in-cloud ventilation by shallow convective clouds and  $\text{CO}_2$ -assimilation is essential to describe the diurnal and vertical distribution of  $\text{CO}_2$  within the lower tropical troposphere.

In the diurnal cycle, the clear-to-cloudy transition allows the formulation of three distinct diurnal regimes depending on the main governing process: entrainment-diluting regime, cloud-ventilation-and-entrainment regime and the  $\text{CO}_2$ -assimilation regime. Shallow convective clouds ( $\sim 23\%$ ), clear air entrainment ( $\sim 21\%$ ), and rainforest  $\text{CO}_2$ -assimilation ( $\sim 56\%$ ) collectively shape the clear-to-cloudy boundary layer  $\text{CO}_2$  budget, with their relative importance varying per diurnal regime. Without dynamically active shallow convective clouds, cloud ventilation is absent, and a stronger entrainment alone governs the diurnal evolution and vertical distribution of  $\text{CO}_2$  within the boundary layer. Since mass removal by clouds is suppressed, the boundary layer can deepen by up to 500 m relative to cloudy conditions. These compensating effects therefore result in a total  $\text{CO}_2$ -flux convergence comparable to a cloudy boundary layer. Nonetheless, entrainment and  $\text{CO}_2$ -assimilation alone are 20%–25% less effective in mixing  $\text{CO}_2$  in the boundary layer, causing the boundary layer to become well-mixed up to an hour later compared to cloudy conditions.

In the vertical, shallow convective clouds actively couple the rainforest surface to the upper atmosphere. Cloud ventilation and entrainment collectively facilitate the vertical exchange of  $\text{CO}_2$  up to heights twice the boundary layer depth and at a rate comparable to  $\text{CO}_2$ -assimilation at the rainforest canopy, significantly affecting its vertical distribution until late afternoon. The clouds organize the turbulent exchange of  $\text{CO}_2$  by extending the positive end of the negative-to-positive  $\text{CO}_2$ - $\text{H}_2\text{O}$  correlation from the boundary layer top upward. Spatial analysis revealed that this turbulent exchange is predominantly confined to cloud scales of  $< 2$  km. Within these dynamic clouds, strong updraughts (fluctuations of vertical velocity up to  $9 \text{ m s}^{-1}$ ) transport  $\text{CO}_2$  and moisture from the boundary layer upward. These cloud driven processes fall within the gray zone of resolving convection, posing a challenge for accurate representation in numerical weather prediction and climatological models. Over time, the vertical negative to positively correlated signal persists, shifting to being primarily driven by fluctuations in  $\text{H}_2\text{O}$  when the lower troposphere becomes well-mixed in  $\text{CO}_2$ .

Although our study is constrained by the available observations, we believe it offers significant insights into the interactions between rainforest  $\text{CO}_2$ -assimilation, clear air entrainment, and in-cloud ventilation by shallow mass fluxes, as well as the organization of turbulent  $\text{CO}_2$  exchange by shallow convective clouds in the lower tropical troposphere. Future research should prioritize assessing the impact of potential misrepresentations of these clear-to-shallow convective processes in numerical weather prediction and climate models, such as the Integrated Forecasting System (IFS) of the European Centre for Medium-Range Weather Forecasts, which underpin most inverse-modeled  $\text{CO}_2$  budget estimates. Greater complexity should also be incorporated by examining the relative contributions of both local and large-scale processes on the turbulent exchange of  $\text{CO}_2$ . Example processes are the role of multi-layer canopy dynamics during the morning transition, large-scale advection, subsidence, the role of shallow-to-deep convective clouds and synoptic weather systems. Ongoing work aims to build upon the framework established in this study, employing detailed observations from campaigns such as CloudRoots-Amazon22 alongside LES and the IFS at resolutions of 4.4, 9, and 25 km, to further investigate and resolve these multi-scale dynamics involved in the turbulent exchange of  $\text{CO}_2$  over the Amazon and ecosystems alike.

## Conflict of Interest

The authors declare no conflicts of interest relevant to this study.

## Data Availability Statement

A copy of the used CloudRoots-Amazon22 data, visualization code, and DALES version 4.4 code are available via the 4-TU repository (de Feiter et al., 2025).

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